

# Critical Exponents of 3D Ising Model from Theory and Monte Carlo Simulations of Very Large Lattices

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# Reorganized perturbation theory

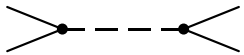
(Ann. Phys. (Leipzig) 10 (2001) 299)

Consider  $\varphi^4$  model with the Hamiltonian

$$H/T = \int (r_0 \varphi^2(\mathbf{x}) + c(\nabla \varphi(\mathbf{x}))^2 + u \varphi^4(\mathbf{x})) d\mathbf{x} ,$$

Grouping of Feynman diagrams  $\Rightarrow$  the Dyson equation

$$\frac{1}{2G_i(\mathbf{k})} = r_0 + ck^2 - \frac{\partial D(G)}{\partial G_i(\mathbf{k})} + \vartheta_i(\mathbf{k})$$

for the correlation function  $\langle \varphi_i(\mathbf{k}) \varphi_j(-\mathbf{k}) \rangle = \delta_{ij} G_i(\mathbf{k})$ . Here  $D(G)$  is the (resummed) sum of grouped *skeleton diagrams* constructed of the fourth order vertices , including *all* original diagrams of  $\varphi^4$  perturbations.

# Advantages

- The method allows to make certain analysis without cutting the series  $\implies$  **exact critical exponents**.
- The asymptotics of  $G(\mathbf{k})$  is found **directly as an expansion in powers of  $k$**  avoiding doubtful intermediate expansions in divergent parameters like  $\ln k$ .

The latter problem is not satisfactory solved in the perturbative RG approach.  $\ln k$  diverges at  $k \rightarrow 0$  and the RG method is not correct, since a contradiction can be derived! – Sec. 2 in **Ann. Phys. (Leipzig) 10 (2001) 299**.

1) correction-to-scaling for  $1/[k^2 G(\mathbf{k})]$  is  $\delta X(\mathbf{k}, \mu) = \mathcal{O}(\epsilon^2)$ , as obtained from a first-principles equation assuming the Wilson–Fisher fixed point; 2) we get  $\delta X(\mathbf{k}, \mu) = \mathcal{O}(\epsilon)$  by matching coefficients at  $\ln k$ , since  $\omega = \epsilon + \mathcal{O}(\epsilon^2)$ .

# Critical exponents

Critical exponents predicted by grouping of Feynman diagrams: for  $n = 1, 2, \dots, d < 4$  (only  $n = 1$  for  $d = 2$ )

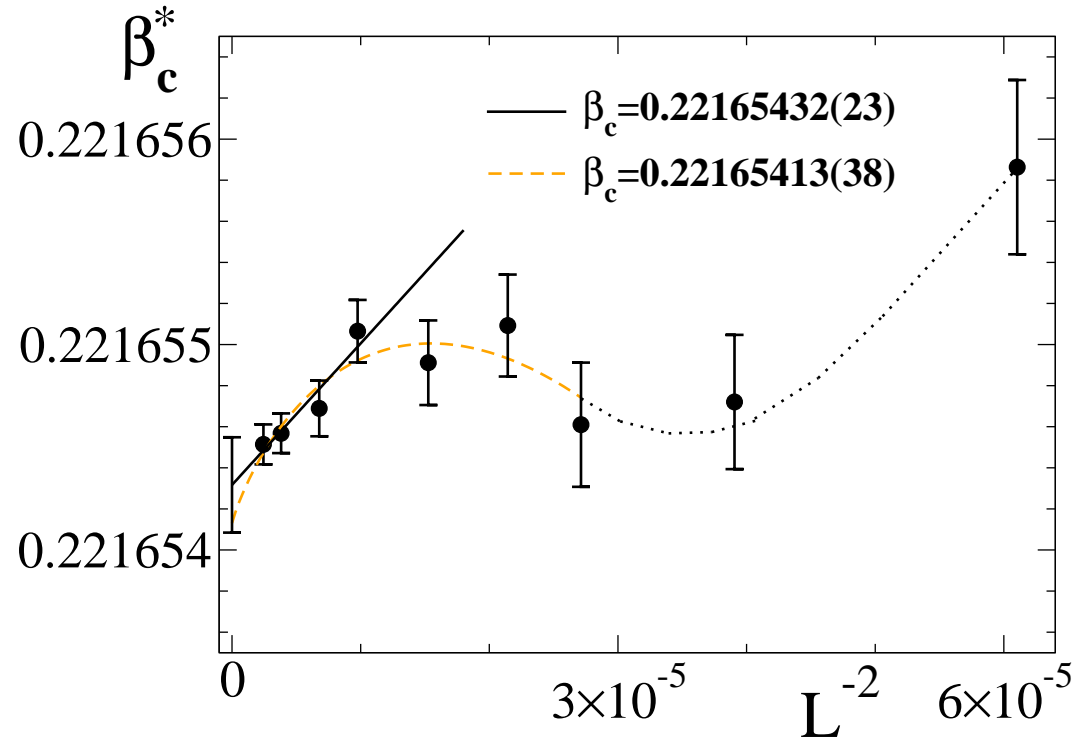
$$\gamma = \frac{d + 2j + 4m}{d(1 + m + j) - 2j} \quad \nu = \frac{2(1 + m) + j}{d(1 + m + j) - 2j},$$

where  $m \geq 1$  and  $j \geq -m$  are integers. It reproduces the known exact critical exponents of 2D Ising model ( $m = 3, j = 0$ ), the mean-field exponents at  $d \rightarrow 4$  (any  $j$  and  $m$ ), as well as those of the spherical model ( $j/m \rightarrow \infty$ ).

**The Ising case:**  $m = 3, j = 0 \Rightarrow \gamma = 7/4, \nu = 1$  at  $d = 2$  and  $\gamma = 5/4, \nu = 2/3$  ( $\alpha = 0, \eta = 1/8, \beta = 3/8$ ) at  $d = 3$ .

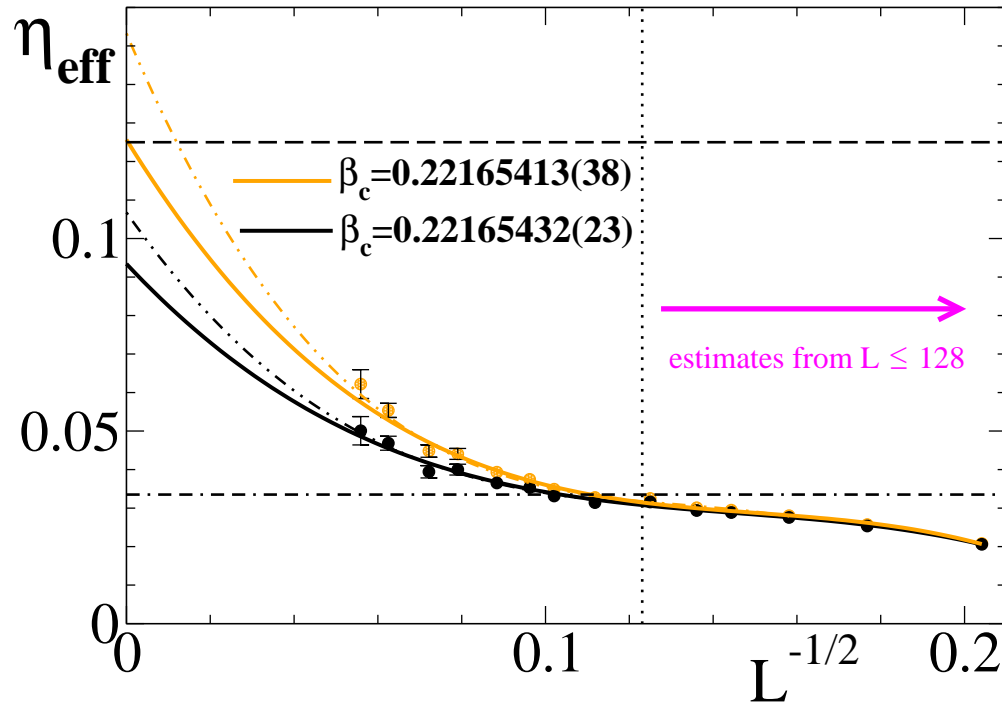
**The conventional (RG) values:**  $\gamma \simeq 1.24, \nu \simeq 0.63, \alpha \simeq 0.11$   
 $\eta \simeq 0.0335, \beta \simeq 0.326$ .

# Estimation of the critical coupling



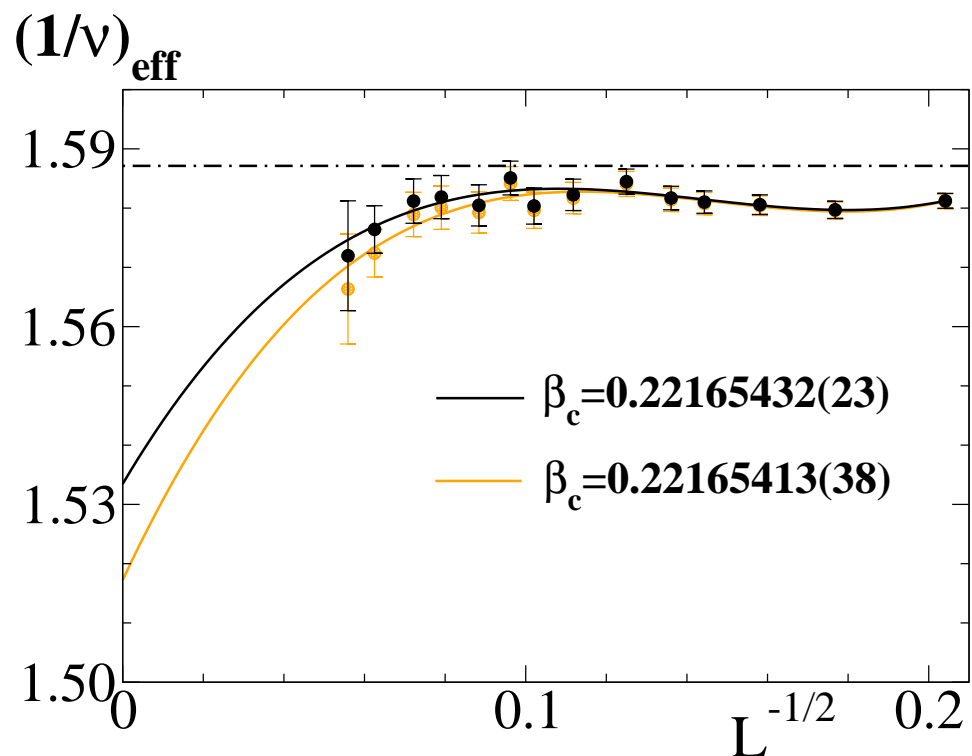
Couplings  $\beta_c^*(L)$  ( $\approx$  Binder cumulant crossing points) obtained from  $\tilde{\beta}_c(L_1)$  data at  $L_1 = L, L/2$  with  $128 \leq L \leq 640$ , where  $\tilde{\beta}_c(L)$  correspond to  $U = \langle m^4 \rangle / \langle m^2 \rangle^2 = 1.6$ . The  $L^{-(1/\nu)-\omega}$  convergence to  $\beta_c$  is expected. Estimation with  $\nu = 2/3$  and  $\omega = 1/2$  (shown) yields  $\beta_c = 0.22165432(23)$  (linear fit) and  $\beta_c = 0.22165413(38)$  (nonlinear fit).  $\nu = 0.63$  and  $\omega = 0.8$  gives  $\beta_c = 0.22165438(18)$  and  $\beta_c = 0.22165429(25)$ , respectively.

# Estimation of $\eta$ in 3D Ising model



The effective critical exponent  $\eta_{\text{eff}}$  vs  $L^{-1/2}$  estimated from susceptibility data at  $\beta = \beta_c$  for each pair of sizes  $(2L; L/2)$  according to  $\chi \sim L^{2-\eta}$ . **The simulated range of sizes [12; 640].** Our theoretical value  $\eta = 1/8$  (dashed line), the RG value  $\eta \simeq 0.0335$  (dot-dashed line).

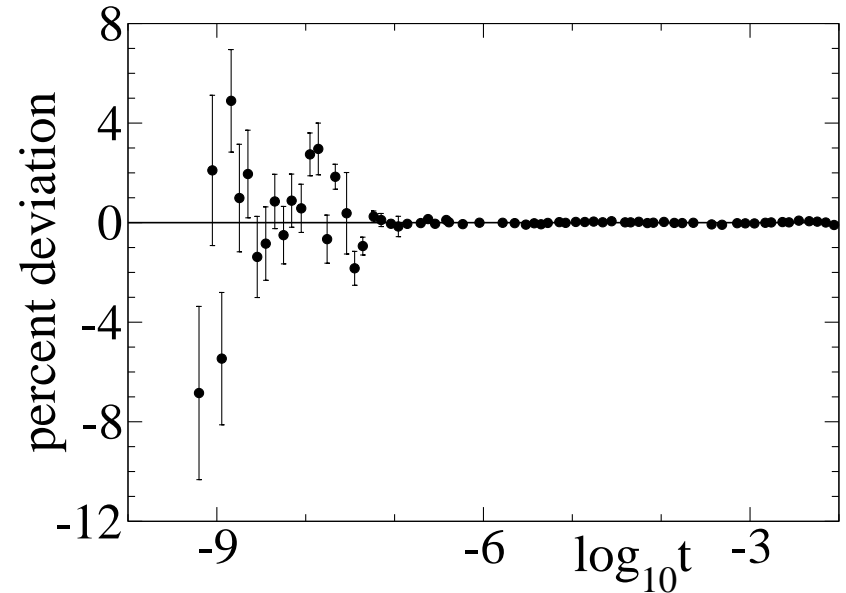
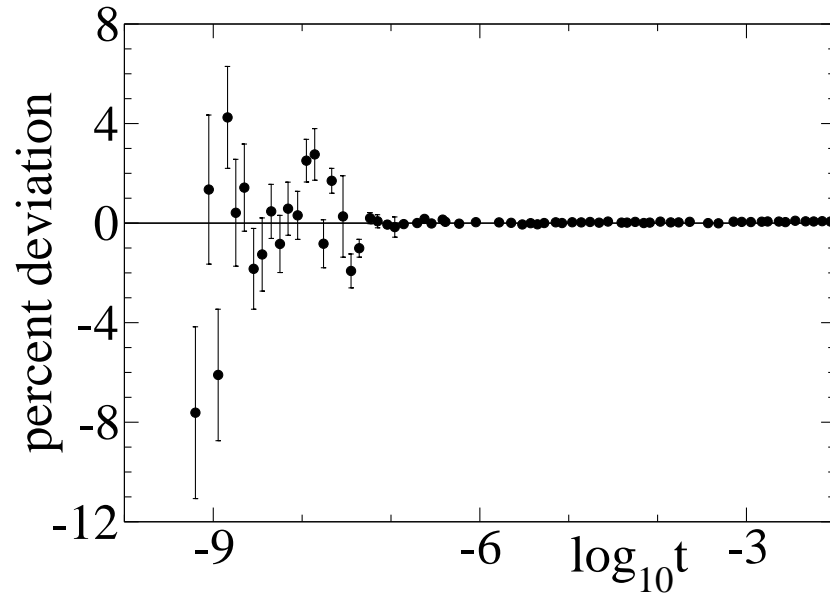
# Estimation of $\nu$ in 3D Ising model



The effective critical exponent  $(1/\nu)_{\text{eff}}$  vs  $L^{-1/2}$  estimated from the derivative of  $\langle m^2 \rangle^2 / \langle m^4 \rangle$  at  $\beta = \beta_c$  (with  $\sim L^{1/\nu}$  scaling) for each pair of sizes  $(2L; L/2)$ . **The simulated range [12; 640]**. Our theoretical value  $1/\nu = 1.5$ , the RG value  $1/\nu \simeq 1.587$  (dot-dashed line).

# Comparison to $C_p$ data in liquid helium

(following experiments by J. A. Lipa, et. al, see [Phys. Rev. B 68, 174518 \(2003\)](#))



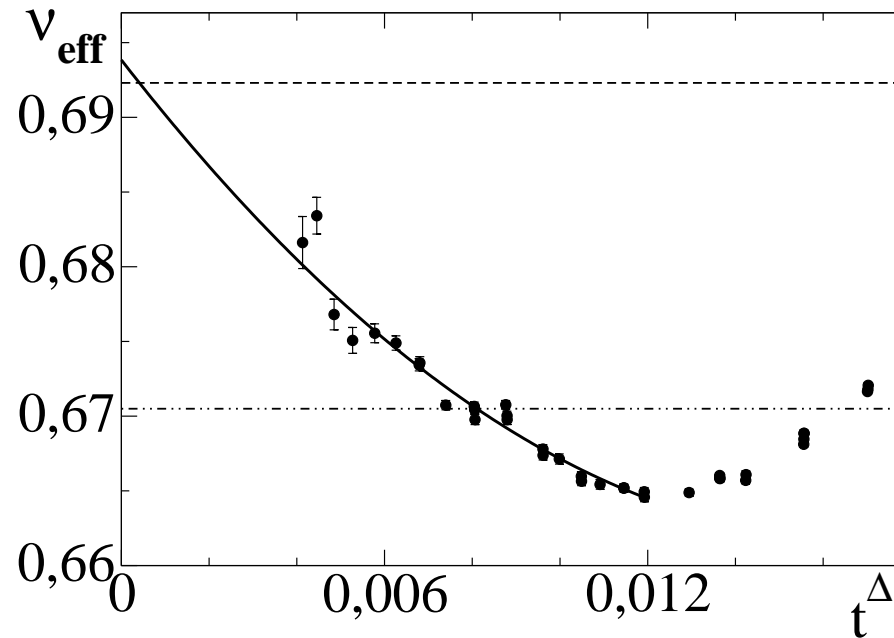
[Eur. Phys. J. B 45, 459 \(2005\)](#): percent deviation from

$$C_p = \frac{A}{\alpha} t^{-\alpha} \left( 1 + at^{\Delta} + bt^{2\Delta} \right) + B, \quad \alpha = -0.0127, \Delta = 0.529 \text{ (left)}$$

$$C_p = t^{-\alpha} (C + A \ln t) \left( 1 + at^{\Delta} \right) + B, \quad \alpha = -1/13, \Delta = 5/13 \text{ (right)}$$



# $\nu$ from superfluid fraction in liquid $He$



**Eur. Phys. J. B 45, 459 (2005)**: The effective exponent  $\nu_{\text{eff}}$  depending on  $t^\Delta$  (with  $\Delta = 5/13$ ) evaluated from local slopes of  $\ln \rho$  vs.  $\ln t$  plot. The lower line – (RG) value 0.6705, the upper line – our theoretical value  $\nu = 9/13$ .

Data taken from **L. S. Goldner, N. Mulders, G. Ahlers, J. Low Temperature Phys. 93, 131 (1993)**.